

Multi-bubble blow-up solutions to stochastic nonlinear Schrödinger equations

Deng Zhang

Shanghai Jiao Tong University

16th Workshop on Markov Processes and Related Topics

Central South University, Changsha

arXiv: 2012.14037

based on the work in joint with Yiming Su

July 15, 2021

- 1 I. Motivations
- 2 II. Main results
- 3 III. Further work

I. Motivations

We are concerned with the **blow-up** of stochastic nonlinear Schrödinger equation:

$$\begin{aligned}dX &= i\Delta X dt + i|X|^{\frac{4}{d}} X dt - \mu X dt + X dW(t), \\ X(0) &= X_0 \in H^1(\mathbb{R}^d).\end{aligned}\tag{SNLS}$$

Here,

I. Motivations

We are concerned with the **blow-up** of stochastic nonlinear Schrödinger equation:

$$\begin{aligned}dX &= i\Delta X dt + i|X|^{\frac{4}{d}}X dt - \mu X dt + X dW(t), \\ X(0) &= X_0 \in H^1(\mathbb{R}^d).\end{aligned}\tag{SNLS}$$

Here,

- **Focusing mass-critical** nonlinearity: **cubic** if $d = 2$, **quintic** if $d = 1$.

I. Motivations

We are concerned with the **blow-up** of stochastic nonlinear Schrödinger equation:

$$\begin{aligned} dX &= i\Delta X dt + i|X|^{\frac{4}{d}} X dt - \mu X dt + X dW(t), \\ X(0) &= X_0 \in H^1(\mathbb{R}^d). \end{aligned} \tag{SNLS}$$

Here,

- **Focusing mass-critical** nonlinearity: **cubic** if $d = 2$, **quintic** if $d = 1$.
- W is the colored Wiener process

$$W(t, x) = \sum_{k=1}^N i\phi_k(x) B_k(t), \quad x \in \mathbb{R}^d, \quad t \geq 0.$$

$\phi_k \in C_b^\infty(\mathbb{R}^d)$, $\operatorname{Re} W = 0$ (*conservative case*),
 B_k are independent real-valued Brownian motions on $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$.

- $X dW$ is taken in the sense of controlled rough path.

I. Motivations

We are concerned with the **blow-up** of stochastic nonlinear Schrödinger equation:

$$\begin{aligned} dX &= i\Delta X dt + i|X|^{\frac{4}{d}} X dt - \mu X dt + X dW(t), \\ X(0) &= X_0 \in H^1(\mathbb{R}^d). \end{aligned} \tag{SNLS}$$

Here,

- **Focusing mass-critical** nonlinearity: **cubic** if $d = 2$, **quintic** if $d = 1$.
- W is the colored Wiener process

$$W(t, x) = \sum_{k=1}^N i\phi_k(x) B_k(t), \quad x \in \mathbb{R}^d, \quad t \geq 0.$$

$\phi_k \in C_b^\infty(\mathbb{R}^d)$, $\operatorname{Re} W = 0$ (*conservative case*),
 B_k are independent real-valued Brownian motions on $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$.

- $X dW$ is taken in the sense of controlled rough path.
- $\mu(x) = \frac{1}{2} \sum_{j=1}^N \phi_k^2(x)$, $x \in \mathbb{R}^d$. (Physic motivation: conservation of mass.)

I. Motivations: SNLS

Physical motivations

- 1 Propagation of nonlinear dispersive waves in random media (e.g. thermal vibration of molecules): $|X(t)|_2^2 = |X_0|_2^2$.
[Bang et al.: Phys. Rev. 1994; Appl. Anal. 1995]

Noise : random potential, approximated by Gaussian noise

- 2 Theory of measurements continuous in time in quantum mechanics.
[Barchielli, Gregoratti: Lecture Notes Physics, 2009]

Noise : stochastic continuous measurement in equation of state X

$\widehat{P}_{X_0}^T(F) := \int_F |X(T)|_2^2 dP$: physical probability of events occurring in $[0, T]$

“Physical Probability Law”, martingale property $(\mathbb{E}|X(t)|_2^2 = |X_0|_2^2)$

I. Motivations: NLS

Nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0, \quad u(0) = u_0 \in H^1(\mathbb{R}^d).$$

(NLS) admits rich invariance and conservation laws:

I. Motivations: NLS

Nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0, \quad u(0) = u_0 \in H^1(\mathbb{R}^d).$$

(NLS) admits rich invariance and conservation laws:

- **Invariance:**

Invariant under the translation, scaling, phase rotation, Galilean transform.

In particular, L^2 -norm is invariant under $u(t, x) \rightarrow \lambda^{-\frac{d}{2}} u(\frac{t}{\lambda^2}, \frac{x}{\lambda}) \Rightarrow$ **mass-critical**.

Pseudo-conformal invariance, particularly important in the mass-critical case, under

$$u(t, x) \rightarrow (-t)^{-\frac{d}{2}} u\left(\frac{1}{-t}, \frac{x}{-t}\right) e^{-i\frac{|x|^2}{4t}}, \quad t \neq 0.$$

I. Motivations: NLS

Nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0, \quad u(0) = u_0 \in H^1(\mathbb{R}^d).$$

(NLS) admits rich invariance and conservation laws:

- **Invariance:**

Invariant under the translation, scaling, phase rotation, Galilean transform.

In particular, L^2 -norm is invariant under $u(t, x) \rightarrow \lambda^{-\frac{d}{2}} u(\frac{t}{\lambda^2}, \frac{x}{\lambda}) \Rightarrow$ **mass-critical**.

Pseudo-conformal invariance, particularly important in the mass-critical case, under

$$u(t, x) \rightarrow (-t)^{-\frac{d}{2}} u\left(\frac{1}{-t}, \frac{x}{-t}\right) e^{-i\frac{|x|^2}{4t}}, \quad t \neq 0.$$

- **Conservation law:**

Mass : $M(u)(t) := \int_{\mathbb{R}^d} |u(t)|^2 dx = M(u_0),$

Energy : $E(u)(t) := \frac{1}{2} \int_{\mathbb{R}^d} |\nabla u(t)|^2 dx - \frac{d}{2d+4} \int_{\mathbb{R}^d} |u(t)|^{2+\frac{4}{d}} dx = E(u_0),$

Momentum : $Mom(u) := \text{Im} \int_{\mathbb{R}^d} \nabla u \bar{u} dx = Mom(u_0).$

I. Motivations: NLS

An important role is played by the **ground state** Q :

which is the unique positive radial solution to the nonlinear elliptic equation

$$\Delta Q - Q + Q^{1+\frac{4}{d}} = 0.$$

Threshold ([Weinstein: CMP 1982]):

The mass of ground state is exactly the **threshold** for GWP and blow-up:

- Solutions exist globally in the subcritical mass case $\|u_0\|_{L^2} < \|Q\|_{L^2}$;
- Singularity can be formed in the (super)critical mass case $\|u_0\|_{L^2} \geq \|Q\|_{L^2}$.

I. Motivations: NLS

Three different blow-up regimes:

Critical mass case $\|u_0\|_{L^2} = \|Q\|_{L^2}$:

I. Motivations: NLS

Three different blow-up regimes:

Critical mass case $\|u_0\|_{L^2} = \|Q\|_{L^2}$:

Two important families of solutions:

- Solitary wave $u(t, x) = Q(x)e^{it}$, exists globally, but does not scatter.

I. Motivations: NLS

Three different blow-up regimes:

Critical mass case $\|u_0\|_{L^2} = \|Q\|_{L^2}$:

Two important families of solutions:

- ① Solitary wave $u(t, x) = Q(x)e^{it}$, exists globally, but does not scatter.
- ② Pseudo-conformal blow-up solution:

$$S_T(t, x) := \frac{1}{(T-t)^{\frac{d}{2}}} Q\left(\frac{x}{T-t}\right) e^{\frac{i}{T-t} - i\frac{|x|^2}{4(T-t)}}, \quad x \in \mathbb{R}^d, \quad t < T.$$

Note:

- S_T is obtained from the solitary wave solution and pseudo-conformal symmetry.
- S_T blows up at time T with rate $(T-t)^{-1}$.
- $\|S_T\|_{L^2} = \|Q\|_{L^2}$

I. Motivations: NLS

Three different blow-up regimes:

Critical mass case $\|u_0\|_{L^2} = \|Q\|_{L^2}$:

Two important families of solutions:

- ① Solitary wave $u(t, x) = Q(x)e^{it}$, exists globally, but does not scatter.
- ② Pseudo-conformal blow-up solution:

$$S_T(t, x) := \frac{1}{(T-t)^{\frac{d}{2}}} Q\left(\frac{x}{T-t}\right) e^{\frac{i}{T-t} - i\frac{|x|^2}{4(T-t)}}, \quad x \in \mathbb{R}^d, \quad t < T.$$

Note:

- S_T is obtained from the solitary wave solution and pseudo-conformal symmetry.
- S_T blows up at time T with rate $(T-t)^{-1}$.
- $\|S_T\|_{L^2} = \|Q\|_{L^2} \Rightarrow$ **minimal mass blow-up solutions**

I. Motivations: NLS

Three different blow-up regimes:

Critical mass case $\|u_0\|_{L^2} = \|Q\|_{L^2}$:

Two important families of solutions:

- ① Solitary wave $u(t, x) = Q(x)e^{it}$, exists globally, but does not scatter.
- ② Pseudo-conformal blow-up solution:

$$S_T(t, x) := \frac{1}{(T-t)^{\frac{d}{2}}} Q\left(\frac{x}{T-t}\right) e^{\frac{i}{T-t} - i\frac{|x|^2}{4(T-t)}}, \quad x \in \mathbb{R}^d, \quad t < T.$$

Note:

- S_T is obtained from the solitary wave solution and pseudo-conformal symmetry.
- S_T blows up at time T with rate $(T-t)^{-1}$.
- $\|S_T\|_{L^2} = \|Q\|_{L^2} \Rightarrow$ **minimal mass blow-up solutions**
- Instability.

I. Motivations: NLS

Rigidity result ([Merle: Duke Math. J. 1993]):

The pseudo-conformal blow-up solution is exactly the unique minimal mass blow-up solution to (NLS), up to the invariance of (NLS).

Construction results for other dispersive equations:

- Mass-critical inhomogeneous Hartree equation [Cao, Su: JMP, 2013],
- Mass-critical gKdV equation [Martel, Merle, Raphaël: JEMS, 2017],
- Modified Benjamin-Ono equation [Martel, Pilod: Math. Ann., 2017].

I. Motivations: NLS

Super-critical mass case with mass slightly above $\|Q\|_{L^2}$: $\|u_0\|_{L^2} = \|Q\|_{L^2} + \varepsilon$

Two different asymptotic scenarios have been observed:

- ① the pseudo-conformal blow-up rate $\|\nabla u\|_{L^2} \sim \frac{1}{T-t}$;
- ② the log-log blow-up rate $\|\nabla u\|_{L^2} \sim \sqrt{\frac{\log|\log(T-t)|}{T-t}}$.

See the series of works by Merle and Raphaël: [Merle, Raphaël: AOM, 2005; GFA, 2003; Invet. Math. 2004, JAMS, 2006].

See also the loglog blow-up solutions with L^2 initial data (low regularity): [Fan, Mendelson: arXiv: 2010.07821].

I. Motivations: NLS

Super-critical mass case with large mass:

The complete characterization of singularity formation is open.

Mass quantization conjecture [Merle, Raphaël: CMP, 2005]:

General blow-up solutions shall concentrate the mass $m_k (\geq \|Q\|_{L^2}^2)$ at the singularities and converge strongly in L^2 to a residue u^* away from the singularities.

I. Motivations: NLS

Super-critical mass case with large mass:

The complete characterization of singularity formation is open.

Mass quantization conjecture [Merle, Raphaël: CMP, 2005]:

General blow-up solutions shall concentrate the mass $m_k (\geq \|Q\|_{L^2}^2)$ at the singularities and converge strongly in L^2 to a residue u^* away from the singularities.

One important step is to construct **multi-bubble blow-up solutions**:

- Multi-bubble blow-up solutions with pseudo-conformal rate [Merle: CMP, 1990];
- Multi-bubble blow-up solutions with loglog rate [Fan: AHP, 2017].

I. Motivations: NLS

Super-critical mass case with large mass:

The complete characterization of singularity formation is open.

Mass quantization conjecture [Merle, Raphaël: CMP, 2005]:

General blow-up solutions shall concentrate the mass $m_k (\geq \|Q\|_{L^2}^2)$ at the singularities and converge strongly in L^2 to a residue u^* away from the singularities.

One important step is to construct **multi-bubble blow-up solutions**:

- Multi-bubble blow-up solutions with pseudo-conformal rate [Merle: CMP, 1990];
- Multi-bubble blow-up solutions with loglog rate [Fan: AIHP, 2017].

Question: What about stochastic nonlinear Schrödinger equation ?

I. Motivations: SNLS

Stochastic case:

- 1 Physical references: mass-critical case ($d = 1, 2$)
[Bang, Christiansen, If, Rasmussen: Phys. Rev. E, 1994; Appl. Anal. 1995],
[Rasmussen, Gaididei, Bang, Christiansen: Phys. Letters A, 1995].
Two dimensional collapse is important in nonlinear optics.
- 2 Numerical study:
[Debussche, Di Menza: Phys. D, 2002; Appl. Math. Lett.],
[Millet, Rodriguez, Roudenko, Yang: arXiv 2005.14266v1, 2006.10695v1].

I. Motivations: SNLS

Stochastic case:

- ① Physical references: mass-critical case ($d = 1, 2$)
 [Bang, Christiansen, If, Rasmussen: Phys. Rev. E, 1994; Appl. Anal. 1995],
 [Rasmussen, Gaididei, Bang, Chrisiansen: Phys. Letters A, 1995].
 Two dimensional collapse is important in nonlinear optics.
- ② Numerical study:
 [Debussche, Di Menza: Phys. D, 2002; Appl. Math. Lett.],
 [Millet, Rodriguez, Roudenko, Yang: arXiv 2005.14266v1, 2006.10695v1].
- ③ Mass-supercritical case:
 [de Bouard, Debussche: PTRF, 2002; AOP, 2005].
 The *conservative* noise has the effect to accelerate blow-up with positive probability.
 The proof is based on a stochastic version of virial evolution initiated in [Glassey: JMP, 1977] and the support theorem.
- ④ Mass-(super)critical case:
 [Barbu, Röckner, Z.: JDE, 2016].
 The *non-conservative* noise has the damped effect to prevent blow-up with high probability.

I. Motivations: SNLS

Stochastic case: Quantitative descriptions of blow-up dynamics

- ① Critical mass case:

[Su, Z.: [arXiv: 2002.09659](#)]

Construction of minimal mass blow-up solutions

I. Motivations: SNLS

Stochastic case: Quantitative descriptions of blow-up dynamics

- ① Critical mass case:

[Su, Z.: [arXiv: 2002.09659](#)]

Construction of minimal mass blow-up solutions \Rightarrow threshold of GWP and blow-up.

I. Motivations: SNLS

Stochastic case: Quantitative descriptions of blow-up dynamics

- 1 Critical mass case:

[Su, Z.: [arXiv: 2002.09659](#)]

Construction of minimal mass blow-up solutions \Rightarrow threshold of GWP and blow-up.

- 2 Supercritical mass case:

[Su, Fan, Z.: [arXiv: 2011.12171](#)]

Construction of log-log blow-up solutions to SNLS.

I. Motivations: SNLS

Stochastic case: Quantitative descriptions of blow-up dynamics

- 1 Critical mass case:

[Su, Z.: [arXiv: 2002.09659](#)]

Construction of minimal mass blow-up solutions \Rightarrow threshold of GWP and blow-up.

- 2 Supercritical mass case:

[Su, Fan, Z.: [arXiv: 2011.12171](#)]

Construction of log-log blow-up solutions to SNLS.

Goal: Study multi-bubble blow-up solutions

- 1 Construction ?
- 2 Uniqueness ?

II. Main results

Assumptions

- Asymptotical flatness (A0): For $1 \leq k \leq N$ and for any multi-index $\nu \neq 0$,

$$\lim_{|x| \rightarrow \infty} \langle x \rangle^2 |\partial_x^\nu \phi_k(x)| = 0.$$

- Flatness at the origin (A1): $\exists \nu_* \in \mathbb{N}$ such that for $1 \leq k \leq N$ and $1 \leq j \leq K$,

$$\partial_x^\nu \phi_k(x_j) = 0, \quad \forall 0 \leq |\nu| \leq \nu_*.$$

Two cases

- Case (I):** $\{x_j\}_{j=1}^K$ are arbitrary distinct points in \mathbb{R}^d , and $\{\omega_j\}_{j=1}^K (\subseteq \mathbb{R}^+)$ satisfy that for some $\omega > 0$, $|\omega_j - \omega| \leq \varepsilon$ for any $1 \leq j \leq K$, where $\varepsilon > 0$;
- Case (II):** $\{\omega_j\}_{j=1}^K$ are arbitrary points in \mathbb{R}^+ , and $\{x_j\}_{j=1}^K (\subseteq \mathbb{R}^d)$ satisfy that $|x_j - x_l| \geq \varepsilon^{-1}$ for any $1 \leq j \neq l \leq K$, where $\varepsilon > 0$.

Example: The singularities $\{x_k\}$ can be arbitrary if $\{\omega_k\}$ are the same.

II. Main results

Theorem (Existence)

Consider $d = 1, 2$. Assume (A0) and (A1) with $\nu_* \geq 5$. Let $K \in \mathbb{N}/\{0\}$, $\{\vartheta_j\} \subseteq \mathbb{R}$, $\{x_j\}_{j=1}^K \subseteq \mathbb{R}^d$ and $\{\omega_j\}_{j=1}^K \subseteq \mathbb{R}^+$ satisfy either Case (I) or Case (II).

Then, for $\mathbb{P} - a.e.$ ω there exists $\varepsilon^*(\omega) > 0$ small enough such that the following holds. For any $\varepsilon \in (0, \varepsilon^*]$, there exists $\tau^* > 0$ small enough such that for any $T \in (0, \tau^*(\omega)]$, there exist $X_0(\omega) \in \Sigma$ and a corresponding blow-up solution $X(\omega)$, satisfying

$$\|e^{-W(t,\omega)}X(t,\omega) - \sum_{j=1}^K S_j(t)\|_{\Sigma} \leq C(T-t)^{\frac{1}{2}(\nu_*-5)+\zeta}, \quad t \in [0, T],$$

where $\Sigma := \{v \in H^1 : |x|v \in L^2\}$, $C > 0$, $\zeta \in (0, 1)$, and S_j , $1 \leq j \leq K$, are the pseudo-conformal blow-up solutions

$$S_j(t, x) = (\omega_j(T-t))^{-\frac{d}{2}} Q\left(\frac{x-x_j}{\omega_j(T-t)}\right) e^{-\frac{i}{4} \frac{|x-x_j|^2}{T-t} + \frac{i}{\omega_j^2(T-t)} + i\vartheta_j}, \quad t \in (0, T).$$

II. Main results

Remark

- The constructed blow-up solution concentrates at the given K points:

$$|X(t, \omega)|^2 \rightarrow \sum_{j=1}^K \|Q\|_{L^2}^2 \delta_{x=x_j}, \quad \text{as } t \rightarrow T.$$

- The asymptotic order can be improved if the noise is more flat at the singularities. If $\nu_* \geq 6$, the asymptotic can be even taken in the more regular space $H^{\frac{3}{2}}$.

II. Main results

Remark

- The constructed blow-up solution concentrates at the given K points:

$$|X(t, \omega)|^2 \rightarrow \sum_{j=1}^K \|Q\|_{L^2}^2 \delta_{x=x_j}, \quad \text{as } t \rightarrow T.$$

- The asymptotic order can be improved if the noise is more flat at the singularities. If $\nu_* \geq 6$, the asymptotic can be even taken in the more regular space $H^{\frac{3}{2}}$.
- Multi-bubble blow-up solutions were first constructed by Merle for NLS when $d \geq 1$, the main blow-up profile is built on any function that decay exponentially, while the frequencies $\{\omega_j\}_{j=1}^K$ are required to have a uniform positive lower bound and the asymptotic is taken in the space $L^{2+\frac{4}{d}}$.

The above multi-bubble solutions are constructed when $d = 1, 2$, the blow-up profile is built on the ground state.

The gain here is that, in Case (I) the frequencies are allowed to be arbitrarily small, and the asymptotic is taken in the energy space H^1 , which is important to study the uniqueness problem.

II. Main results

Theorem (Conditional uniqueness)

Consider $d = 1, 2$. Assume (A0) and (A1) with $\nu_* \geq 5$. Let $K \in \mathbb{N}$, $\{\vartheta_j\} \subseteq \mathbb{R}$, $\{x_j\}_{j=1}^K \subseteq \mathbb{R}^d$ and $\{\omega_j\}_{j=1}^K \subseteq \mathbb{R}^+$, satisfy Case (I) or Case (II).

Then, for $\mathbb{P} - a.e.$ ω there exists $\varepsilon^*(\omega) > 0$ very small such that the following holds. For any $\varepsilon \in (0, \varepsilon^*]$, there exists $\tau^* > 0$ small enough such that for any $T \in (0, \tau^*(\omega)]$, there exists a unique blow-up solution $X(\omega)$ satisfying

$$\|e^{-W(t,\omega)}X(t,\omega) - \sum_{j=1}^K S_j(t)\|_{H^1} \leq C(T-t)^{3+\zeta}, \quad t \in [0, T),$$

where $C > 0$, $\zeta \in (0, 1)$, and $\{S_j\}$ are the pseudo-conformal blow-up solutions.

II. Main results

Application 1: minimal mass blow-up solutions

Theorem

Consider $d = 1, 2$. Assume (A0) and (A1), $\nu_* \geq 5$. Let $x_*, \omega_*, \vartheta_*$ be any given points.

Then, for \mathbb{P} -a.e. ω there exists $\tau^*(\omega) > 0$ sufficiently small, such that for any $T \in (0, \tau^*(\omega)]$ there exists a minimal mass blow-up solution $X(\omega)$ satisfying

$$\|e^{-W(t,\omega)}X(t,\omega) - S_*(t)\|_{\Sigma} \leq C(T-t)^{\frac{1}{2}(\nu_*-5)+\zeta}, \quad t \in [0, T),$$

where $C > 0$, $\zeta \in (0, 1)$, S_* is as above with $x_*, \omega_*, \vartheta_*$ replacing $\alpha_j, \omega_j, \vartheta_j$, respectively.

Moreover, in the case where $\nu_* \geq 11$, there exists a unique minimal mass blow-up solution $X(\omega)$ satisfying

$$\|e^{-W(t,\omega)}X(t,\omega) - S_*(t)\|_{H^1} \leq C(T-t)^{3+\zeta}, \quad t \in [0, T).$$

II. Main results

Application 2: NLS with lower perturbations

$$i\partial_t v + \Delta v + |v|^{\frac{4}{d}} v + a_1 \cdot \nabla u + a_2 u = 0,$$

where

$$a_1(t, x) = 2i \sum_{k=1}^N \nabla \phi_k(x) h_k(t),$$

$$a_2(t, x) = - \sum_{j=1}^d \left(\sum_{k=1}^N \partial_j \phi_k(x) h_k(t) \right)^2 + i \sum_{k=1}^N \Delta \phi_k(x) h_k(t),$$

ϕ_k satisfies Assumptions (A0) and (A1) and $h_k \in C(\mathbb{R}^+; \mathbb{R})$, $1 \leq k \leq K$.

In particular, if $h_k = B_k$, then the above equation is called RNLS.

These results are also applicable to NLS, in which $a_1, a_2 \equiv 0$.

Note: Absence of the pseudo-conformal invariance and the conservation law of energy.

III. Further work

Further work: [Cao, Su, Zhang: arXiv: 2105.14554].

Consider the canonical focusing L^2 -critical nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0.$$

Uniqueness Problem: does there exist a **unique** multi-bubble blow-up solution v to (NLS) such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{H^1} = o(1), \quad \text{as } t \rightarrow T. \quad (o(1) \text{ Condition})$$

III. Further work

Further work: [Cao, Su, Zhang: arXiv: 2105.14554].

Consider the canonical focusing L^2 -critical nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0.$$

Uniqueness Problem: does there exist a **unique** multi-bubble blow-up solution v to (NLS) such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{H^1} = o(1), \quad \text{as } t \rightarrow T. \quad (o(1) \text{ Condition})$$

Remark:

- Related to the open problem of uniqueness of **multi-solitons** [Martel: ICM 2018].

III. Further work

Further work: [Cao, Su, Zhang: arXiv: 2105.14554].

Consider the canonical focusing L^2 -critical nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0.$$

Uniqueness Problem: does there exist a **unique** multi-bubble blow-up solution v to (NLS) such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{H^1} = o(1), \quad \text{as } t \rightarrow T. \quad (o(1) \text{ Condition})$$

Remark:

- Related to the open problem of uniqueness of **multi-solitons** [Martel: ICM 2018].
- Related to the solitary wave conjecture.

III. Further work

Further work: [Cao, Su, Zhang: arXiv: 2105.14554].

Consider the canonical focusing L^2 -critical nonlinear Schrödinger equation (NLS):

$$i\partial_t u + \Delta u + |u|^{\frac{4}{d}} u = 0.$$

Uniqueness Problem: does there exist a **unique** multi-bubble blow-up solution v to (NLS) such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{H^1} = o(1), \quad \text{as } t \rightarrow T. \quad (o(1) \text{ Condition})$$

Remark:

- Related to the open problem of uniqueness of **multi-solitons** [Martel: ICM 2018].
- Related to the solitary wave conjecture.
- Uniqueness is known in the single bubble case by the seminal paper [Merle: Duke Math. J., 1993]. The strong rigidity result was also proved by [Raphaël, Szeftel: JAMS, 2011] for inhomogeneous NLS.
- Uniqueness is open in the multi-bubble case. It provides an illustration of the rigidity of equation (NLS) around the pseudo-conformal blow-up solutions.

III. Further work

Remark:

- Related to the *mass quantization conjecture* in [Merle, Raphaël: CMP, 2005]: it states that blow-up solutions to (NLS) shall concentrate the mass $m_k (\geq \|Q\|_{L^2}^2)$ at the singularities and converge strongly in L^2 to a residue u^* away from the singularities. The above uniqueness problem provides one uniqueness class for such solutions.

III. Further work

Remark:

- Related to the *mass quantization conjecture* in [Merle, Raphaël: CMP, 2005]: it states that blow-up solutions to (NLS) shall concentrate the mass $m_k (\geq \|Q\|_{L^2}^2)$ at the singularities and converge strongly in L^2 to a residue u^* away from the singularities. The above uniqueness problem provides one uniqueness class for such solutions.
- Similarity with local uniqueness problem of bubbling solutions to nonlinear elliptic equations. [Cao, Peng, Yan: Cambridge University Press, 2021]

For instance, the local uniqueness of multi-peak solutions u_ε to Brezis-Nirenberg problem concentrating at different points $\{x_1, \dots, x_K\}$ satisfying

$$\|u_\varepsilon - \sum_{k=1}^K P U_{x_{k,\varepsilon}, \lambda_{k,\varepsilon}}\|_{H^1} = o(1), \quad x_{k,\varepsilon} \rightarrow x_k, \text{ and } \lambda_{k,\varepsilon} \rightarrow +\infty, \text{ as } \varepsilon \rightarrow 0,$$

where P is the projection from $H^1(\Omega)$ onto $H_0^1(\Omega)$, Ω is a smooth and bounded domain in \mathbb{R}^d , and $U_{x,\lambda}$ solves the elliptic equation $\Delta u + u \frac{d+2}{d-2} = 0$ in \mathbb{R}^d , $d \geq 5$.

III. Further work

Theorem (Multi-bubble blow-up case)

Consider (NLS) in dimensions $d = 1, 2$. Let $T \in \mathbb{R}$, $K \in \mathbb{N} \setminus \{0\}$. Let $\{\vartheta_k\} \subseteq \mathbb{R}$, $\{\omega_k\}$ and $\{x_k\}$ satisfy either Case (I) or Case (II).

Then, for any $\zeta \in (0, 1)$, there exists $\varepsilon^* > 0$ such that the following holds. For any $0 < \varepsilon < \varepsilon^*$, there exists a unique multi-bubble blow-up solution v such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{L^2} + (T - t) \|\nabla v(t) - \sum_{k=1}^K \nabla S_k(t)\|_{L^2} = o(1), \quad \text{as } t \text{ close to } T,$$

and additionally double average condition

$$\frac{1}{T-t} \int_t^T \frac{1}{T-s} \int_s^T \|v(r) - \sum_{k=1}^K S_k(r)\|_{H^1}^2 dr ds = \mathcal{O}((T-t)^\zeta), \quad (0+) - \text{Condition}$$

where S_k , $1 \leq k \leq K$, are the pseudo-conformal blow-up solutions.

III. Further work

Theorem (continued)

Moreover, the unique solution v converges exponentially fast to $\sum_{k=1}^K S_k$ in the pseudo-conformal space, i.e., there exists $\delta > 0$ such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{\Sigma} = \mathcal{O}(e^{-\frac{\delta}{T-t}}), \quad \text{for } t \text{ close to } T.$$

In particular, the above results hold for the multi-bubble blow-up solutions v such that

$$\|v(t) - \sum_{k=1}^K S_k(t)\|_{H^1} = \mathcal{O}((T-t)^\zeta), \quad \text{for } t \text{ close to } T. \quad (0+ \text{ Condition})$$

Thank you